COSMOLOGICAL HYDROGEN RECOMBINATION:
THE EFFECT OF HIGH-N STATES

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OUTLINE

- Cosmological Recombination in a nutshell

- Breaking the naive model

- Why should you care? Effects on CMB, inferences about primordial physics

- Our tools

- Preliminary results!

# COSMOLOGICAL RECOMBINATION IN A NUTSHELL

- $p + e^- \leftrightarrow H^{(n)} + \gamma^{(nc)}$ outrun by Hubble expansion, Saha Eq. fails: **Relic free electrons!**

- For continuum $\to 1s$, $2p \to 1s$, $t_{\mathrm{mf}} < H^{-1}(T)$

- Two ways out of $n = 2$ bottleneck

  - Two-photon decays: $H^{2s} \to H^{1s} + \gamma + \gamma$ $\quad \boxed{\mathbf{\Lambda_{2s \to 1s} = 8.22 \ s^{-1}}}$

  - Redshifting out of resonance: $\boxed{R \sim \lambda_\alpha^{-3} (\dot{a}/a) n_{1s}^{-1}}$

- Peebles (1967) assumes radiative eq. between n=2 and excited states, and between eq. between *angular momentum (l)* substates

$$\boxed{\mathcal{N}_n = \mathcal{N}_2 e^{-(E_n - E_2)/T}}$$

$$\boxed{\mathcal{N}_{nl} = \mathcal{N}_n \frac{(2l+1)}{n^2}}$$

PEEBLES MODEL ASSUMPTIONS/RESULTS

- Blackbody spectrum of radiation assumed
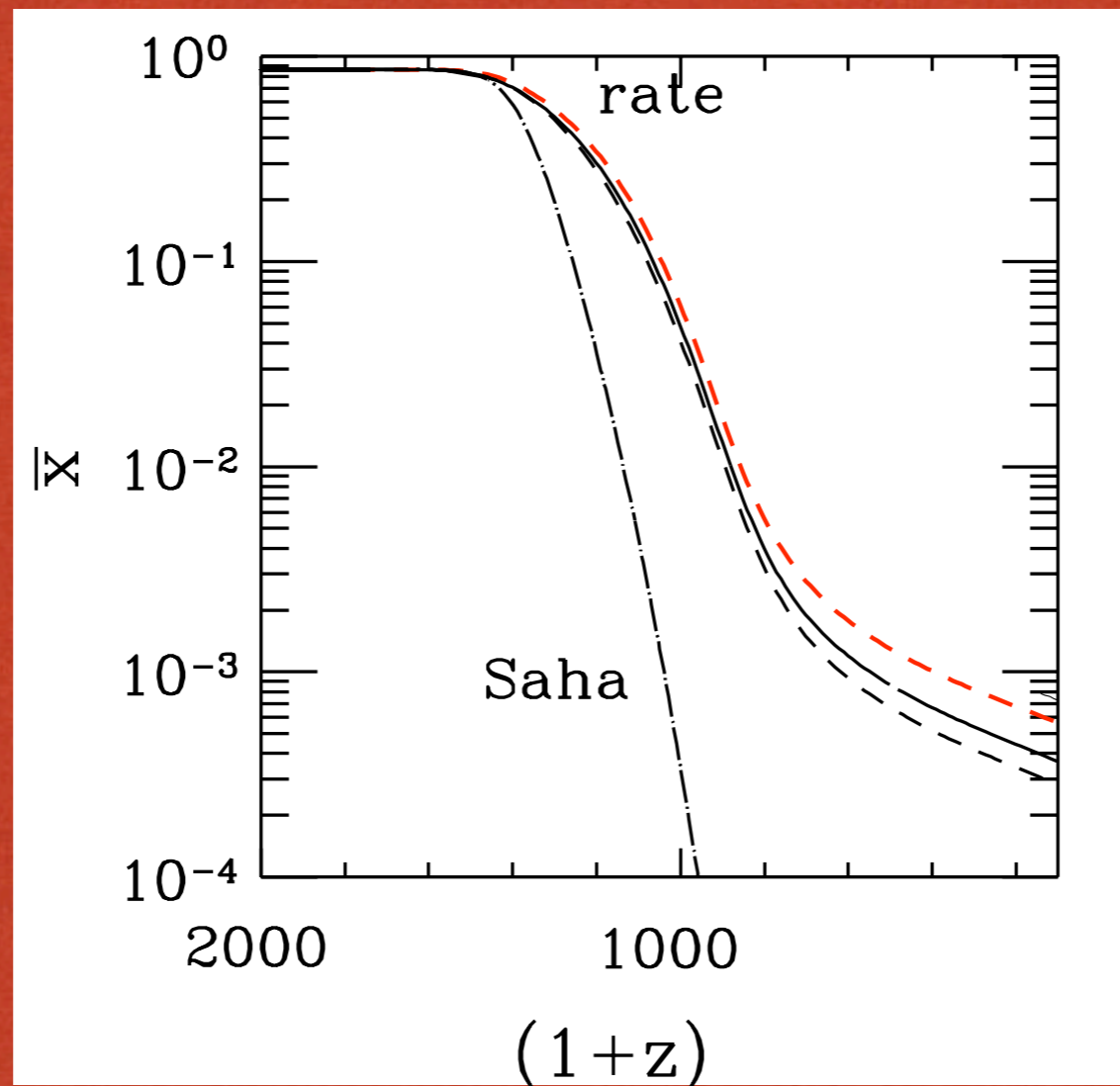
- State of the Art for 30 years!

BREAKING THE NAIVE MODEL

- ~~Equilibrium~~ between *l states* (z<900)

- Cool radiation field: ~~Boltzmann eq. of higher n~~ (z<900)

- Radiation field is non-thermal

- Compton scattering inefficient $T_m \neq T_\gamma$ (z<500)

- Higher level 2γ transitions (Hirata 2008) (z<800)

- Lyman-α diffusion

BREAKING THE NAIVE MODEL

- Radiation field is cool: ~~Boltzmann eq. of higher n~~

- Treated by Seager et al. (2000) $n_{\max} = 300$ RecFAST!!!

- ~~Equilibrium~~ between *l states*

- Treated by Chluba et al. (2005) for $n_{\max} = 100$

- Beyond this, testing convergence with $n_{\max}$ is hard!

$$t_{\mathrm{compute}} \sim \mathcal{O}\left(\mathrm{weeks}\right)$$

How to proceed if we want 0.1% accuracy in $x_e(z)$?

BREAKING THE NAIVE MODEL

- Radiation field is cool: ~~Boltzmann eq. of higher n~~

- Treated by Seager et al. (2000) $n_{\max} = 300$ RecFAST!!!

- ~~Eq.~~ between *l states*: dipole selection bottleneck: $\Delta l = \pm 1$

- Treated by Chluba et al. (2005) for $n_{\max} = 100$

- Beyond this, testing convergence with $n_{\max}$ is hard!
  $$t_{\mathrm{compute}} \sim \mathcal{O}\,(\mathrm{weeks})$$

WHY PROCEED?

- $\gamma - e^-$ decouple at $z_{\rm dec} \simeq 1090$, during recombination

- If $\Delta z_{\rm rec} \uparrow$, γ have more time to diffusively (Silk) damp C_l

- If $\Delta z_{\rm rec} \uparrow$, low-$\lambda$ modes cancel more along l.o.s.

- Polarization $\Pi \propto \Theta_2(\eta)$, $\Theta_2(\eta)$ until $g \sim 1$, so $\Delta z_{\rm rec} \uparrow \rightarrow \Pi \uparrow$

- Modes with $\boxed{\lambda < H^{-1}(\eta)}$ suppressed : $\boxed{C_l \rightarrow C_l e^{-2\tau}}$

- Planck will be CV limited out to $l \sim 2500$:
  Possible degeneracies between
 cosmo parameters (e.g. n_s) and atomic physics!

THE MULTI-LEVEL ATOM

- Bound-free rates (Hummer and Storey 1965)

- Bound-bound rates (Hummer and Storey 1965)

- Solve for escape fraction of lines with Sobolev approx

- Add line photons and follow evolution of $f(E_\gamma, T)$

- Excited state populations obey $\dfrac{d\vec{x}}{dt} = \mathbf{R}\vec{x} + \vec{s}$

- Rate on LHS is $\boxed{t_{\text{rec}}^{-1} \sim 10^{-12} s^{-1} \ll \mathbf{R} \, , \vec{s} \to \vec{x} \simeq \mathbf{R}^{-1}\vec{s}}$

  - Steady state approximation good for n>1.

- Matrix is $\sim n_{max}^2 \times n_{max}^2$

- Brute force would require for $n_{max}^6 \sim 1000$ s for $n_{max} = 200$ for a single time step

- Sparsity to the rescue $\boxed{\mathbf{M}_{l,l-1}\vec{x}_{l-1} + \mathbf{M}_{l,l}\vec{x}_l + \mathbf{M}_{l,l+1}\vec{x}_{l+1} = \vec{s}_l}$
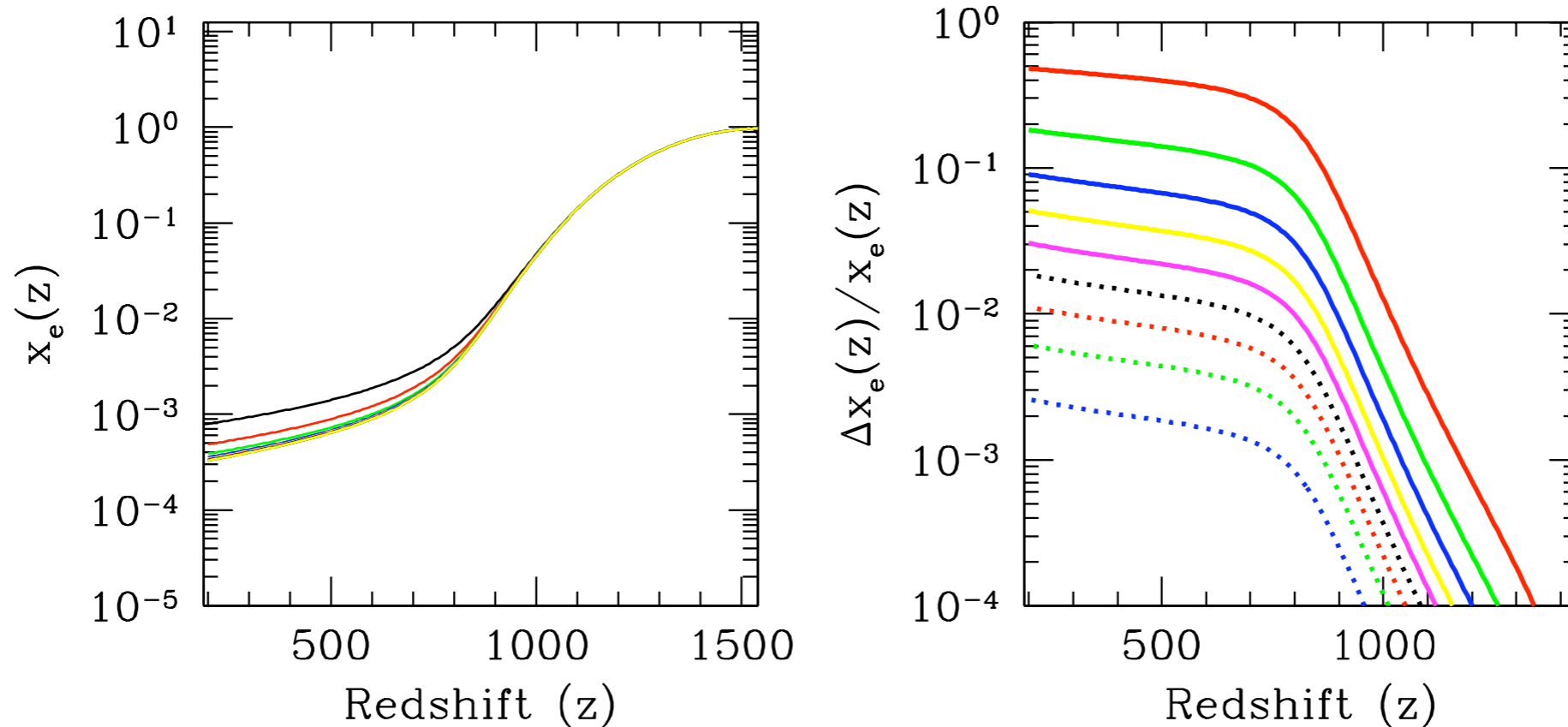
- $x_e(z)$ falls with increasing $n_{max} = 10 \rightarrow 100$, as expected:

- Rec Rate>downward BB Rate> Ionization, upward BB rate

- Even for $n_{max} = 100$, code computes in only 2 hours

- To add: Spectral distortions, collisional rates, increase n_{max}